Magnetic properties of nanostructured Ca_{1-x}Gd_xMnO₃ obtained by glycine-nitrate procedure

V. Spasojevic, V. Kusigerski, M. Rosic, J. Blanusa, M. Perovic, A. Mrakovic, B. Antic, and B. Matovic

The Vinca Institute, University of Belgrade, P.O. Box 522, Belgrade, Serbia vojas@vinca.rs

Ca_{1-x}Gd_xMnO₃ nanopowders (x=0.05, 0.10, 0.15, 0.20) with perovskite type crystal structure were synthesized by the glycine nitrate procedure [1]. Starting powders were prepared by combustion of solutions containing mixture of glycine with metal nitrates in their appropriate stoichiometric ratios. The so-obtained powders were annealed at the temperature of 850 °C for 10 minutes to produce final nanostructured samples with the average nanoparticle size of about 20 nm. Magnetic measurements show that electron doping by Gd^{3+} ions substantially changes $CaMnO_3$ antiferromagnetic (AFM) behaviour. After introduction of Gd^{3+} ions, significant ferromagnetic (FM) component appears due to an emergence of double exchange interaction between Mn3+-Mn4+ ions. This resulted in appearance of a low temperature plateau in field cooled (FC) magnetization as well as in emergence of hysteresis loop with the relatively high coercivity up to 2300 Oe. Presence of competing long-range AFM and shortrange FM interactions and their randomness lead to a frustration of manganese magnetic moments, and appearance of the spin-glass state at low temperature of about T_{SG}≈65 K (Fig.1, Inset). Concentration dependence of effective magnetic moments $\mu_{eff}(x)$ is linear, while Curie-Weiss temperature $\theta(x)$ changes its sign from negative to positive for concentrations above x=0.05 (Fig. 2) which points to the formation of FM clusters in AFM matrix [2]. The same picture is supported by low temperature M(H) measurements were maximums in H_c(x) and M(x) curves at certain concentration x is a consequence of the balance between FM-cluster/AF-matrix interactions (Fig. 3). Recorded large field-cooled hysteresis shift H_{shift} (Fig.4.) can be also considered as a consequence of exchange bias that emerges at the borderline of FM-AFM regions. We assume that these interactions are also responsible for the unusual $H_c(T)$ plateau at the temperatures below spin-glass transitions (Fig. 4).

References

- [1] D. Markovic, V. Kusigerski, M. Tadic, J. Blanusa, V. Antisari, V. Spasojevic, Scr. Mater. **59** (2008) 35.
- [2] I. Sudheendra, A. R. Raju and C. N. Rao, J. Phys.: Condens. Matter 15 (2003) 895.

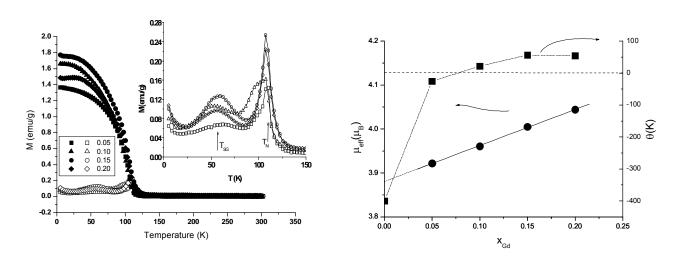
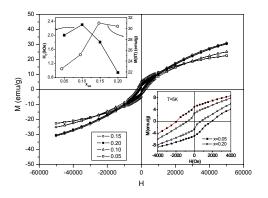


Fig. 1. Fig. 2.



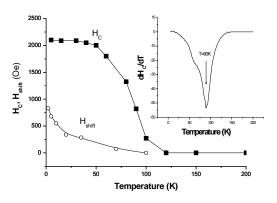


Fig. 3. Fig. 4.